## Torque on Current Loop Immersed in Constant Magnetic Field

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**Setup** Consider a closed loop of current I lying on a plane, immersed within a constant magnetic field  $\vec{B}$ . We will now discuss why the total magnetic force on this loop is zero; but the total torque  $\vec{\tau}$  is non-zero and in fact is the cross product between the current loop's magnetic moment  $\vec{\mu}$  and the magnetic field:

$$\vec{\tau} = \vec{\mu} \times \vec{B}.\tag{0.0.1}$$

The magnetic moment is, in turn, defined as the directed area of the loop times its current:

$$\vec{\mu} = I\vec{A} = IA\hat{n},\tag{0.0.2}$$

where  $\hat{n}$  is the unit vector perpendicular to the loop given by the 'right hand rule'. Namely, if the  $\hat{n}$  is pointed at you while you are looking down at the loop, the current needs to appear flowing in the counter-clockwise direction.

**Total Force** Remember the force on an infinitesimal segment of a current carrying wire is

$$\mathrm{d}\vec{F} = I\mathrm{d}\vec{\ell} \times \vec{B}.\tag{0.0.3}$$

The total force is therefore

$$\vec{F} = I\left(\oint d\vec{\ell}\right) \times \vec{B} \tag{0.0.4}$$

$$= I\left(\oint \mathrm{d}x\widehat{x} + \oint \mathrm{d}y\widehat{y} + \oint \mathrm{d}z\widehat{z}\right) \times \vec{B}$$
(0.0.5)

$$= I \left[ x \hat{x} + y \hat{y} + z \hat{z} \right]_{\text{initial position on loop}}^{\text{final position on loop}} \times \vec{B} = \vec{0}.$$
(0.0.6)

The initial and final positions are the same – we are integrating around a closed loop – and hence the integral is zero.

Area Before tackling the torque we first note that the directed area  $\vec{A}$  can itself be expressed as

$$\vec{A} = A\hat{n} = \frac{1}{2} \oint \vec{r} \times d\vec{\ell}.$$
(0.0.7)

The  $\vec{r} = (x, y, z)$  is the position vector joining the (arbitrary) origin of the coordinate system to the point (x, y, z) on the current loop; whereas  $d\vec{\ell} = \hat{x}dx + \hat{y}dy + \hat{z}dz$  in turn is the infinitesimal

displacement on the loop along the direction of the current. To understand eq. (0.0.7), we note that the infinitesimal triangle formed by the sides  $\vec{r}, \vec{r} + d\vec{\ell}$  and  $d\vec{\ell}$  has area

$$dA = \frac{1}{2} |\vec{r} + d\vec{\ell}| |d\vec{\ell}| \sin \psi \approx \frac{1}{2} |\vec{r}| |d\vec{\ell}| \sin \psi, \qquad (0.0.8)$$

where  $\psi$  is the angle between  $d\vec{l}$  and either  $\vec{r}$  or  $\vec{r} + d\vec{l}$  (the difference between the two is infinitesimal). But remember  $|\vec{a} \times \vec{b}| = |\vec{a}||\vec{b}|\sin\theta$ , where  $\theta$  is the angle between  $\vec{a}$  and  $\vec{b}$ . We have arrived at

$$\mathrm{d}A = \frac{1}{2} |\vec{r} \times \mathrm{d}\vec{\ell}|. \tag{0.0.9}$$

The directed area is therefore eq. (0.0.7). Finally, the magnetic moment of a single current loop is

$$\vec{\mu} = I\vec{A} = \frac{I}{2} \oint \vec{r} \times d\vec{\ell}.$$
(0.0.10)

**Torque** The torque exerted upon an infinitesimal segment of current carrying wire is

$$d\vec{\tau} = \vec{r} \times d\vec{F} = \vec{r} \times \left( I d\vec{\ell} \times \vec{B} \right), \qquad (0.0.11)$$

where we have used eq. (0.0.3). Employing the identity

$$\vec{a} \times \left(\vec{b} \times \vec{c}\right) = (\vec{a} \cdot \vec{c})\vec{b} - (\vec{a} \cdot \vec{b})\vec{c} \tag{0.0.12}$$

then yields

$$\vec{\tau} = I \oint \left( (\vec{r} \cdot \vec{B}) d\vec{\ell} - (\vec{r} \cdot d\vec{\ell}) \vec{B} \right). \tag{0.0.13}$$

Note that the second term involves

$$\oint \vec{r} \cdot d\vec{\ell} = \oint (xdx + ydy + zdz) = \oint \frac{1}{2} \left( d(x^2) + d(y^2) + d(z^2) \right)$$
(0.0.14)

$$= \frac{1}{2} \left[ x^2 + y^2 + z^2 \right]_{\text{initial point}}^{\text{final point}} = 0.$$
 (0.0.15)

Therefore, the total torque is

$$\vec{\tau} = I \oint (\vec{r} \cdot \vec{B}) \mathrm{d}\vec{\ell}. \tag{0.0.16}$$

Note that the definition of total torque is actually independent of the choice of origin. That is, suppose we choose a different origin, so  $\vec{r'} \equiv \vec{r} + \vec{d}$  (for constant  $\vec{d}$ ):

$$\vec{\tau} = \oint \left( \vec{r} + \vec{d} \right) \times \left( I \mathrm{d}\vec{\ell} \times \vec{B} \right) \tag{0.0.17}$$

$$=\oint \left(\vec{r} \times \left(I \mathrm{d}\vec{\ell} \times \vec{B}\right) + \vec{d} \times \left(I \mathrm{d}\vec{\ell} \times \vec{B}\right)\right). \tag{0.0.18}$$

In the second term, only the  $d\vec{\ell}$  is involved in the integration, since  $\vec{d}$  and  $\vec{B}$  are constant; but  $\oint d\vec{\ell} = [x\hat{x} + y\hat{y} + z\hat{z}]_{\text{initial}}^{\text{final}} = \vec{0}$  and we have

$$\vec{\tau} = \oint \vec{r'} \times \left( I \mathrm{d}\vec{\ell} \times \vec{B} \right) = \oint \vec{r} \times \left( I \mathrm{d}\vec{\ell} \times \vec{B} \right). \tag{0.0.19}$$

We turn to

$$\vec{\mu} \times \vec{B} = \frac{I}{2} \left( \oint \vec{r} \times d\vec{\ell} \right) \times \vec{B} = -\frac{I}{2} \vec{B} \times \left( \oint \vec{r} \times d\vec{\ell} \right) \tag{0.0.20}$$

$$= -\frac{I}{2} \left( \oint \vec{r} (\vec{B} \cdot d\vec{\ell}) - \oint (\vec{B} \cdot \vec{r}) d\vec{\ell} \right).$$
(0.0.21)

Let us examine the first term:

$$\oint \vec{r}(\vec{B} \cdot d\vec{\ell}) = \oint (x, y, z) (B_x dx + B_y dy + B_z dz).$$
(0.0.22)

We may integrate by parts – remember  $\int_a^b f dg = [fg]_a^b - \int_a^b g df$ , and since we have a closed loop the 'surface terms'  $[\ldots]_a^b$  actually vanish – to deduce

$$\oint \vec{r}(\vec{B} \cdot d\vec{\ell}) = -\oint (dx, dy, dz)(B_x x + B_y y + B_z z)$$
(0.0.23)

$$= -\oint (\vec{B} \cdot \vec{r}) \mathrm{d}\vec{\ell}. \tag{0.0.24}$$

In other words, the two terms in eq. (0.0.21) are equal and we have

$$\vec{\mu} \times \vec{B} = I \oint (\vec{B} \cdot \vec{r}) \mathrm{d}\vec{\ell}. \tag{0.0.25}$$

Comparing equations (0.0.25) and (0.0.16), we have arrived at eq. (0.0.1).